

Optimal Control of Dissipative Solids via Vanishing Viscosity

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V E R S I T Ä T

DFG Deutsche
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SPP 1962

Non-smooth and Complementarity-based
Distributed Parameter Systems:
Simulation and Hierarchical Optimization

What are rate-independent systems and
what are vanishing viscosity solutions?

Rate-independent systems

- Rate-independent evolutions are driven by external forces on a time scale much slower than their internal scale.
- When the external loadings undergo a time rescaling, the solutions to the rescaled system are the rescaled solutions of the original system.
- examples: dry friction, plasticity, fracture



A. Mielke, T. Roubířek. **Rate-Independent Systems**. Springer, 2015.

Rate-independent systems

For an initial value $z_0 \in \mathcal{Z}$ find a curve $z : [0, T] \rightarrow \mathcal{Z}$ such that

$$0 \in \partial\mathcal{R}(\dot{z}(t)) + D_z\mathcal{I}(\ell(t), z(t)) \text{ f.a.a. } t \in [0, T], \quad z(0) = z_0.$$

- $z \in \mathcal{Z}$ *state variable in the state space*
- $\mathcal{I} : \mathcal{Z}^* \times \mathcal{Z} \rightarrow \mathbb{R}, \mathcal{I}(\ell, z) = \mathcal{I}_0(z) - \langle \ell, z \rangle$ *potential energy*
- $\ell \in W^{1,\infty}(0, T; \mathcal{Z}^*)$ *external load*
- $\mathcal{R} : \mathcal{Z} \rightarrow [0, \infty)$ convex, l.s.c., **pos. 1-homogeneous** *dissipation potential*
(e.g.: a norm on \mathcal{Z})
- $\partial\mathcal{R} : \mathcal{Z} \rightrightarrows \mathcal{Z}^*$ *convex subdifferential*
- $D_z\mathcal{I}(\ell, \cdot) : \mathcal{Z} \rightarrow \mathcal{Z}^*$ *Fréchet-differential of $z \mapsto \mathcal{I}(\ell, z)$*

Rate-independent systems

For an initial value $z_0 \in \mathcal{Z}$ find a curve $z : [0, T] \rightarrow \mathcal{Z}$ such that

$$0 \in \partial \mathcal{R}(\dot{z}(t)) + D_z \mathcal{I}(\ell(t), z(t)) \text{ f.a.a. } t \in [0, T], \quad z(0) = z_0.$$

interpretations:

- force balance in \mathcal{Z}^* :
dissipative force $\partial \mathcal{R}(\dot{z})$ versus potential restoring force $-D_z \mathcal{I}(\ell, z)$
- if $\mathcal{R}(w) = \frac{1}{2} \|w\|_{\mathcal{Z}}^2$, we would have to solve the gradient flow

$$\dot{z}(t) = -D_z \mathcal{I}(\ell(t), z(t))$$

Vanishing viscosity solutions

problem: for non-convex \mathcal{I} , solutions may have jumps

strategy: ► approximate the system by adding artificial viscosity:

assume $\mathcal{Z} \in \mathcal{V} \subset \mathcal{X}$ and $\mathcal{R}(z) = C\|z\|_{\mathcal{X}}$; for $\varepsilon > 0$, set

$$\mathcal{R}_\varepsilon(z) := \mathcal{R}(z) + \frac{\varepsilon}{2}\|z\|_{\mathcal{V}}^2$$

and solve

$$0 \in \partial \mathcal{R}_\varepsilon(\dot{z}_\varepsilon(t)) + D_z \mathcal{I}(\ell(t), z_\varepsilon(t)) \text{ a.e.}, \quad z_\varepsilon(0) = z_0$$

- 1 obtain a sequence of continuous solutions $(z_\varepsilon)_{\varepsilon>0}$
- 2 show existence of a converging subsequence for $\varepsilon \rightarrow 0$
- 3 find out in which way the limit solves the original problem

- 1 Rate-independent systems and vanishing viscosity solutions
- 2 Existence of viscous solutions
- 3 Vanishing viscosity analysis

Existence of viscous solutions - assumptions

$\mathcal{I}(\ell, z) := \frac{1}{2} \langle Az, z \rangle_{\mathcal{Z}} + \mathcal{F}(z) - \langle \ell, z \rangle_{\mathcal{V}}$ for $z \in \mathcal{Z}$, where

- $A : \mathcal{Z} \rightarrow \mathcal{Z}^*$ is linear, bounded, self-adjoint and coercive
- $\mathcal{F} \in C^2(\mathcal{Z}, \mathbb{R}_{>0})$ fulfills
 - ▶ $D_z \mathcal{F} : \mathcal{Z} \rightarrow \mathcal{V}^*$ is weakly continuous and $D_z \mathcal{F} \in C^1(\mathcal{Z}, \mathcal{V}^*)$,
 - ▶ $\exists C > 0, q \geq 1 \forall z, v \in \mathcal{Z} : \|D_z^2 \mathcal{F}(z)v\|_{\mathcal{V}^*} \leq C(1 + \|z\|_{\mathcal{Z}}^q) \|v\|_{\mathcal{Z}}$

Proposition (Existence of viscous solutions)

There is a solution $z_\varepsilon \in W^{1,1}([0, T], \mathcal{V})$ of the viscous system, fulfilling the energy dissipation balance:

$$\int_s^t \mathcal{R}_\varepsilon(\dot{z}_\varepsilon(r)) + \mathcal{R}_\varepsilon^*(-D_z \mathcal{I}(\ell(r), z_\varepsilon(r))) \, dr + \mathcal{I}(\ell(t), z_\varepsilon(t)) \\ = \mathcal{I}(\ell(s), z_\varepsilon(s)) + \int_s^t \partial_r \mathcal{I}(\ell(r), z_\varepsilon(r)) \, dr \quad \text{for all } 0 \leq s \leq t \leq T.$$

Proof.



Mielke, Rossi, Savaré. **Nonsmooth analysis of doubly nonlinear evolution equations.** 2013.

[▶ show details](#)

There is a constant $C > 0$ such that for all $\varepsilon > 0$ and all solutions z_ε , it holds:

$$\|z_\varepsilon\|_{L^\infty(0, T; \mathcal{Z})} + \sqrt{\varepsilon} \|z_\varepsilon\|_{H^1(0, T; \mathcal{V})} \leq C.$$

There is a sequence of solutions fulfilling

$$\sup_{\varepsilon > 0} \left(\|z_\varepsilon\|_{W^{1,1}(0, T; \mathcal{Z})} + \varepsilon \|z_\varepsilon\|_{W^{1,\infty}(0, T; \mathcal{V})} \right) \leq C.$$

- 1 Rate-independent systems and vanishing viscosity solutions
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Vanishing viscosity analysis

goal: solve differential inclusion for $\varepsilon = 0$ by passing to the limit of the viscous solutions z_ε

problem:

- 1 a priori estimates in non-reflexive spaces
- 2 limit element is not differentiable in general

strategy:

- 1 rescale solutions such that rescaled solutions \hat{z}_ε are uniformly equicontinuous, use Arzelà-Ascoli and obtain a limit \hat{z}
- 2 introduce generalized notion of differentiability fulfilled by \hat{z}

Vanishing viscosity analysis: time parameterization

recall the energy dissipation balance:

$$\begin{aligned} & \mathcal{I}(\ell(s), z_\varepsilon(s)) - \mathcal{I}(\ell(t), z_\varepsilon(t)) + \int_s^t \partial_r \mathcal{I}(\ell(r), z_\varepsilon(r)) \, dr \\ &= \int_s^t \mathcal{R}(\dot{z}_\varepsilon(r)) + \underbrace{\frac{\varepsilon}{2} \|\dot{z}_\varepsilon(s)\|_{\mathcal{V}}^2 + \frac{1}{2\varepsilon} \text{dist}_{\mathcal{V}^*}(-D_z \mathcal{I}(\ell(r), z_\varepsilon(r)), \partial \mathcal{R}(0))^2}_{\xrightarrow{\varepsilon \rightarrow 0} ?} \, dr \end{aligned}$$

Vanishing viscosity analysis: time parameterization

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$$\text{set } \begin{cases} \mathfrak{p}(v, w) & := \inf_{\varepsilon > 0} \left(\mathcal{R}(v) + \frac{\varepsilon}{2} \|v\|_{\mathcal{V}}^2 + \frac{1}{2\varepsilon} \text{dist}_{\mathcal{V}^*}(w, \partial \mathcal{R}(0))^2 \right) \\ s_\varepsilon(t) & := t + \int_0^t \mathfrak{p}(\dot{z}_\varepsilon(r), -D_z \mathcal{I}(\ell(r), z_\varepsilon(r))) dr, \\ S_\varepsilon & := s_\varepsilon(T). \end{cases}$$

Since s_ε is strictly increasing, the inverse $\hat{t}_\varepsilon := s_\varepsilon^{-1} : [0, S_\varepsilon] \rightarrow [0, T]$ exists, and we define

$$\hat{z}_\varepsilon : [0, S_\varepsilon] \rightarrow \mathcal{Z}, \quad \hat{z}_\varepsilon(s) := z_\varepsilon(\hat{t}_\varepsilon(s)).$$

Vanishing viscosity analysis: time parameterization

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$$\text{set } \begin{cases} \mathfrak{p}(v, w) & := \mathcal{R}(v) + \|v\|_{\mathcal{V}} \text{dist}_{\mathcal{V}^*}(w, \partial \mathcal{R}(0)) \text{ for } v \in \mathcal{Z}, w \in \mathcal{V}^* \\ s_\varepsilon(t) & := t + \int_0^t \mathfrak{p}(\dot{z}_\varepsilon(r), -D_z \mathcal{I}(\ell(r), z_\varepsilon(r))) \, dr, \\ S_\varepsilon & := s_\varepsilon(T). \end{cases}$$

Since s_ε is strictly increasing, the inverse $\hat{t}_\varepsilon := s_\varepsilon^{-1} : [0, S_\varepsilon] \rightarrow [0, T]$ exists, and we define

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Vanishing viscosity analysis: Arzelà-Ascoli argument

- thanks to the choice of parameterization, it holds almost everywhere

$$\dot{\hat{t}}_\varepsilon(s) + \mathcal{R}(\dot{\hat{z}}_\varepsilon(s)) + \|\dot{\hat{z}}_\varepsilon(s)\|_{\mathcal{V}} \text{dist}_{\mathcal{V}^*}(-D_z \mathcal{I}(\ell(\hat{t}_\varepsilon(s)), \hat{z}_\varepsilon(s)), \partial \mathcal{R}(0)) = 1$$

▶ therefore, it holds that $[0, S] := \bigcup_\varepsilon [0, S_\varepsilon]$ is compact

- Ehrling's lemma yields $W^{1,\infty}(0, S; \mathcal{X}) \cap L^\infty(0, S; \mathcal{Z}) \subset C([0, S], \mathcal{V})$

▶ thus, $\sup_\varepsilon \mathcal{R}(\dot{\hat{z}}_\varepsilon(s)) \leq 1$ implies uniform equicontinuity of \hat{z}_ε [▶ details](#)

- a priori estimate $\sup_\varepsilon \|z_\varepsilon\|_{L^\infty(0, T; \mathcal{Z})} < \infty$ implies that the set $K := \{\hat{z}_\varepsilon(s) \mid \varepsilon > 0, s \in [0, S]\}$ is compact w.r.t. $\xrightarrow{\mathcal{Z}}$.

▶ thus, there is $\hat{z} \in C([0, S], \mathcal{V})$ such that $\hat{z}_{\varepsilon_n}(s) \xrightarrow{n \rightarrow \infty} \hat{z}(s)$ in \mathcal{Z}

Vanishing viscosity analysis: \mathcal{R} -absolute continuity

Definition and Proposition

A map $f : [0, T] \rightarrow \mathbb{R}$ is called **absolutely continuous** if and only if

$$\exists g \in L^1(0, T) : \forall 0 \leq r < s \leq T : |f(s) - f(r)| \leq \int_r^s g(t) dt.$$

In this case, for almost all $t \in [0, T]$, the limit

$$f'(t) := \lim_{h \rightarrow 0} \frac{f(t+h) - f(t)}{h}$$

exists and $|f'|$ is the least function fulfilling the above estimate.

Vanishing viscosity analysis: \mathcal{R} -absolute continuity

Definition and Proposition

A map $z : [0, T] \rightarrow \mathcal{Z}$ is called **\mathcal{R} -absolutely continuous** if and only if

$$\exists g \in L^1(0, T) : \forall 0 \leq r < s \leq T : \mathcal{R}(z(s) - z(r)) \leq \int_r^s g(t) dt.$$

In this case, for almost all $t \in [0, T]$, the limit

$$\mathcal{R}[z'](t) := \lim_{h \searrow 0} \frac{\mathcal{R}(z(t+h) - z(t))}{h} \in \mathbb{R}$$

exists and $\mathcal{R}[z']$ is the least function fulfilling the above estimate.

$\mathcal{R}[z']$ is called the **generalized metric derivative**.



Ambrosio, Gigli, Savaré. **Gradient Flows in Metric Spaces and in the Space of Probability Measures**. 2008.

Vanishing viscosity analysis: regularity

From the identity

$$\dot{\hat{t}}_\varepsilon(s) + \underbrace{\mathcal{R}(\dot{\hat{z}}_\varepsilon(s))}_{\hat{=}\mathcal{R}[\dot{\hat{z}}'_\varepsilon](s)} + \|\dot{\hat{z}}_\varepsilon(s)\|_{\mathcal{V}} \underbrace{\text{dist}_{\mathcal{V}^*}(-D_z \mathcal{I}(\ell(\hat{t}_\varepsilon(s)), \hat{z}_\varepsilon(s)), \partial \mathcal{R}(0))}_{=: \mathfrak{e}(\hat{t}_\varepsilon(s), \hat{z}_\varepsilon(s))} = 1,$$

we conclude that

- there is $\hat{t} \in W^{1,\infty}(0, S)$ such that $\hat{t}_\varepsilon \xrightarrow{*} \hat{t}$,
- $\hat{z} \in \text{AC}([0, S]; \mathcal{R})$ with $\mathcal{R}[\hat{z}'] \leq 1$,
- for $[a, b] \subset \{s \in [0, S] \mid \mathfrak{e}(\hat{t}(s), \hat{z}(s)) > 0\}$, it holds $\hat{z} \in \text{AC}([a, b]; \mathcal{V})$.



Mielke, Rossi, Savaré. **Balanced viscosity solutions to infinite-dimensional rate-independent systems**. 2016.

Energy dissipation balance - upper bound

Change of variable in the energy dissipation balance yields

$$\begin{aligned} & \mathcal{I}(\ell(\hat{t}_\varepsilon(0)), \hat{z}_\varepsilon(0)) + \int_0^s \partial_\ell \mathcal{I}(\ell(\hat{t}_\varepsilon(r)), \hat{z}_\varepsilon(r)) \dot{\ell}(\hat{t}_\varepsilon(r)) \hat{t}_\varepsilon(r) \, dr \\ &= \mathcal{I}(\ell(\hat{t}(s)), \hat{z}_\varepsilon(s)) + \int_0^s \mathcal{R}(\dot{\hat{z}}_\varepsilon(r)) \, dr \\ &+ \int_0^s \frac{\varepsilon}{2\hat{t}_\varepsilon(r)} \|\dot{\hat{z}}_\varepsilon(r)\|_{\mathcal{V}}^2 + \frac{\hat{t}_\varepsilon(r)}{2\varepsilon} \text{dist}_{\mathcal{V}^*}(-D_z \mathcal{I}(\ell(\hat{t}_\varepsilon(r)), \hat{z}_\varepsilon(r)), \partial \mathcal{R}(0))^2 \, dr. \end{aligned}$$

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By lower semicontinuity arguments, we obtain the estimate

$$\begin{aligned} & \mathcal{I}(\ell(\hat{t}(0)), \hat{z}(0)) + \int_0^s \partial_\ell \mathcal{I}(\ell(\hat{t}(r)), \hat{z}(r)) \dot{\ell}(\hat{t}(r)) \dot{\hat{t}}(r) \, dr \\ & \geq \mathcal{I}(\ell(\hat{t}(s)), \hat{z}(s)) + \int_0^s \mathcal{R}[\dot{\hat{z}}](r) + \|\dot{\hat{z}}(r)\|_{\mathcal{V}} \mathfrak{e}(\hat{t}(r), \hat{z}(r)) \, dr. \end{aligned}$$

Vanishing viscosity solutions - existence

Definition and Theorem

There are $(t, z) \in W^{1,\infty}(0, S) \times AC(0, S; \mathcal{R})$ fulfilling:

$$t(S) = T, \quad \|z\|_{L^\infty(0, S; \mathcal{Z})} < \infty,$$

$$f.a.a. s \in [0, S] : \begin{cases} \dot{t}(s) \geq 0, & \dot{t}(s)\epsilon(t(s), z(s)) = 0, \\ \dot{t}(s) + \mathcal{R}[z'](s) + \|\dot{z}(s)\|_{\mathcal{V}}\epsilon(t(s), z(s)) = 1, \end{cases}$$

for $G := \{s \in [0, S] \mid \epsilon(t(s), z(s)) > 0\}$, we have $z \in AC_{loc}(G; \mathcal{V})$,

$$\begin{aligned} \mathcal{I}(\ell(0), z_0) + \int_0^S \partial_\ell \mathcal{I}(\ell(t(r)), z(r)) \dot{\ell}(t(r)) \dot{t}(r) dr \\ = \mathcal{I}(\ell(t(s)), z(s)) + \int_0^S \mathcal{R}[z'](r) + \|\dot{z}(r)\|_{\mathcal{V}} \epsilon(t(r), z(r)) dr \end{aligned}$$

A triple $(S, t, z) \in [T, \infty) \times W^{1,\infty}(0, S) \times AC(0, S; \mathcal{R})$ fulfilling the above conditions is called **normalized parameterized solution** of the system $0 \in \partial \mathcal{R}(\dot{z}(t)) + D_z \mathcal{I}(\ell(t), z(t))$ a.e. in $[0, T]$, $z(0) = z_0$.

Vanishing viscosity solutions - interpretation

We have the complementarity condition

$$\dot{t}(s) \operatorname{dist}_{\mathcal{V}^*}(-D_z \mathcal{I}(\ell(t(s)), z(s)), \partial \mathcal{R}(0)) = 0.$$

- If $\dot{t}(s) > 0$, then $\operatorname{dist}_{\mathcal{V}^*}(-D_z \mathcal{I}(\ell(t(s)), z(s)), \partial \mathcal{R}(0)) = 0$ and $z(s)$ is a stationary solution of the rate-independent system.
 - ▶ $\dot{z}(s) = 0$: sticking
 - ▶ $\dot{z}(s) \neq 0$: rate-independent sliding
- If $\operatorname{dist}_{\mathcal{V}^*}(-D_z \mathcal{I}(\ell(t(s)), z(s)), \partial \mathcal{R}(0)) > 0$, then $\dot{t}(s) = 0$, and due to the normalization condition, we have $\dot{z}(s) \neq 0$. Thus, the artificial time is frozen and the state changes according to viscous laws. In the outer time scale, this is seen as a jump.

Summary

- ▶ Rate-independent systems and vanishing viscosity solutions
 - approximate system by adding artificial viscosity
- ▶ Solve the viscous system
- ▶ Find limit for vanishing viscosity
 - find suitable parameterization
 - introduce metric derivatives
 - existence and regularity of limit element
 - identify energy-dissipation balance
- ▶ Next goal: optimal control with control variable ℓ
 - prove sufficient a priori estimates/compactness results

Existence of viscous solutions via time discretization

- choose a partition $\{0 = t_0^\tau < t_1^\tau < \dots < t_{N-1}^\tau < T \leq t_N^\tau\}$ of $[0, T]$ with step length $\tau > 0$

- set $z_0^{\tau, \varepsilon} := z_0$ and obtain $(z_k^{\tau, \varepsilon})_{k=1, \dots, N}$ as solutions of

$$z_{k+1}^{\tau, \varepsilon} \in \operatorname{Argmin} \left\{ \mathcal{I}(\ell(t_{k+1}^\tau), z) + \tau \mathcal{R}_\varepsilon \left(\frac{z - z_k^{\tau, \varepsilon}}{\tau} \right), z \in \mathcal{Z} \right\}.$$

- interpolants $\hat{z}_\varepsilon^\tau, \bar{z}_\varepsilon^\tau$ fulfill the energy dissipation estimate

$$\begin{aligned} \int_0^T \mathcal{R}_\varepsilon(\dot{\hat{z}}_\varepsilon^\tau(t)) + \mathcal{R}_\varepsilon^*(-D_z \mathcal{I}(\ell(\underline{\mathbf{t}}^\tau(t)), \bar{z}_\varepsilon^\tau(t))) dt + \mathcal{I}(\ell(T), \hat{z}_\varepsilon^\tau(T)) \\ \leq \mathcal{I}(\ell(0), z_0) + \int_0^T \partial_t \mathcal{I}(\ell(t), \hat{z}_\varepsilon^\tau(t)) dt + \int_0^T r_\tau(t) dt \end{aligned}$$

- this carries over to the limit due to the lower semicontinuity of $\mathcal{R}, \mathcal{R}^*$

Existence of viscous solutions via time discretization

- we have the chain rule

$$\frac{d}{dt}\mathcal{I}(\ell(t), z_\varepsilon(t)) = \langle D_z\mathcal{I}(\ell(t), z_\varepsilon(t)), \dot{z}_\varepsilon(t) \rangle_{\mathcal{Z}} + \partial_t\mathcal{I}(\ell(t), z_\varepsilon(t))$$

- Fenchel-Moreau inequality

$$\mathcal{R}_\varepsilon(\dot{z}_\varepsilon(r)) + \mathcal{R}_\varepsilon^*(-D_z\mathcal{I}(\ell(r), z_\varepsilon(r))) \geq \langle -D_z\mathcal{I}(\ell(r), z_\varepsilon(r)), \dot{z}_\varepsilon(r) \rangle_{\mathcal{Z}}$$

and integration w.r.t. time then yield the opposite estimate

$$\begin{aligned} \int_s^t \mathcal{R}_\varepsilon(\dot{z}_\varepsilon(r)) + \mathcal{R}_\varepsilon^*(-D_z\mathcal{I}(\ell(r), z_\varepsilon(r))) \, dr + \mathcal{I}(\ell(t), z_\varepsilon(t)) \\ \geq \mathcal{I}(\ell(s), z_\varepsilon(s)) + \int_s^t \partial_t\mathcal{I}(\ell(r), z_\varepsilon(r)) \, dr \end{aligned}$$

Existence of viscous solutions via time discretization

- we have the chain rule

$$\begin{aligned} \int_s^t \langle -D_z \mathcal{I}(\ell(r), z_\varepsilon(r)), \dot{z}_\varepsilon(r) \rangle_{\mathcal{Z}} dr + \mathcal{I}(\ell(t), z_\varepsilon(t)) \\ = \mathcal{I}(\ell(s), z_\varepsilon(s)) + \int_s^t \partial_t \mathcal{I}(\ell(r), z_\varepsilon(r)) dr \end{aligned}$$

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and integration w.r.t. time then yield the opposite estimate

$$\begin{aligned} \int_s^t \mathcal{R}_\varepsilon(\dot{z}_\varepsilon(r)) + \mathcal{R}_\varepsilon^*(-D_z \mathcal{I}(\ell(r), z_\varepsilon(r))) dr + \mathcal{I}(\ell(t), z_\varepsilon(t)) \\ \geq \mathcal{I}(\ell(s), z_\varepsilon(s)) + \int_s^t \partial_t \mathcal{I}(\ell(r), z_\varepsilon(r)) dr \end{aligned}$$

Energy dissipation balance - lower bound

For curves (t, z) of the given regularity, $s \mapsto \mathcal{I}(\ell(t(s)), z(s))$ is absolutely continuous and it holds almost everywhere

$$\begin{aligned} \left| \frac{d}{ds} \mathcal{I}(\ell(t(s)), z(s)) - \partial_t \mathcal{I}(\ell(t(s)), z(s)) \dot{t}(s) \right| \\ \leq \mathcal{R}[z'](s) + \|\dot{z}(s)\|_{\mathcal{V}} \epsilon(t(s), z(s)). \end{aligned}$$

Thus, integration w.r.t. s yields the opposite estimate and hence the following identity

$$\begin{aligned} \mathcal{I}(\ell(\hat{t}(0)), \hat{z}(0)) + \int_0^s \partial_\ell \mathcal{I}(\ell(\hat{t}(r)), \hat{z}(r)) \dot{\ell}(\hat{t}(r)) \dot{\hat{t}}(r) dr \\ = \mathcal{I}(\ell(\hat{t}(s)), \hat{z}(s)) + \int_0^s \mathcal{R}[\hat{z}'](r) + \|\dot{\hat{z}}(r)\|_{\mathcal{V}} \epsilon(\hat{t}(r), \hat{z}(r)) dr. \end{aligned}$$





Vanishing viscosity solutions - a priori bounds

Lemma

For all $L > 0$, there is a constant $C_L > 0$ such that for all $z_0 \in \mathcal{Z}$, $\ell \in W^{1,\infty}(0, T; \mathcal{V}^)$ fulfilling $\|z_0\|_{\mathcal{Z}} + \|\ell\|_{H^1(0, T; \mathcal{V}^*)} \leq L$ and all normalized parameterized solutions (S, t, z) associated with the initial value z_0 and the external load ℓ , it holds:*

$$S + \|z\|_{L^\infty(0, S; \mathcal{Z})} + \mathcal{I}(\ell(t(S)), z(S)) < C_L.$$

Literature

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